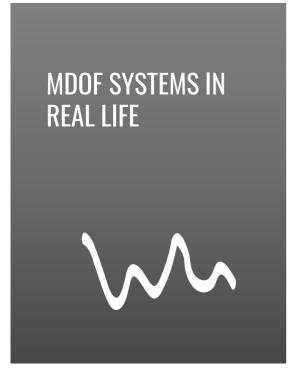
Multiple DOF systems

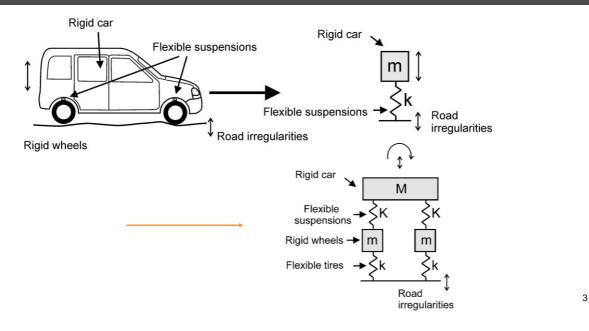




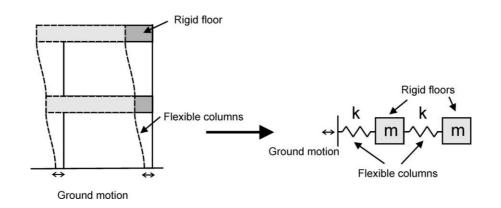




From SDOF to MDOF

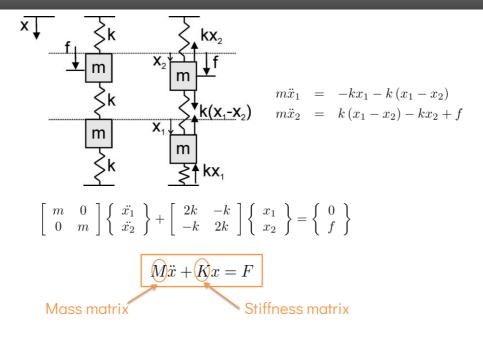


MDOF systems in real life





Equations of motion



6

Free response

$$\begin{array}{c}
M\ddot{x} + Kx = 0 \\
\downarrow \\
(K + r^2 M) \psi = 0
\end{array} \qquad \left\{ \begin{array}{c}
x_1(t) \\
x_2(t)
\end{array} \right\} = \left\{ \begin{array}{c}
A_1 \\
A_2
\end{array} \right\} e^{rt} = \psi e^{rt}$$

Admits a non trivial solution if

$$det(K+r^2M) = 0$$

 r^2 is negative (K and M are positive definite matrices)

$$r^2 = -\omega^2$$

$$\longrightarrow (K - \omega^2 M) \psi = 0$$

7

Free response

$$\left(K - \omega^2 M\right)\psi = 0$$

Generalized eigenvalue problem (- ω^2)

 $det(K - \omega^2 M) = 0$

If the system has *n* degrees of freedom, there exist *n* values of $-\omega^2$ for which this equation is satisfied. These are the *n* eigenvalues which correspond to <u>*n* eigenfrequencies</u>

n eigen vectors ψ are associated to these eigenfrequencies. They correspond to the <u>*n* mode shapes</u> of the structure

The general solution is written in the form: n

$$x(t) = \sum_{i=1}^{n} \left(Z_{i1} cos(\omega_i t) + Z_{i2} sin(\omega_i t) \right) \psi_i$$

8

7

Mode shapes orthogonality

$$\psi_i^T M \psi_j = \delta_{ij} \mu_i$$

$$\psi_i^T K \psi_j = \delta_{ij} \mu_i \omega_i^2$$

Proof :

$$(K - \omega_i^2 M) \psi_i = 0$$

$$(K - \omega_j^2 M) \psi_j = 0$$

$$\omega_i \neq \omega_j$$

$$(1)$$

$$(2)$$

Premultiply (1) by ψ_j^T , (2) by ψ_i^T and substract taking into account symmetry of K ($\psi_i^T K \psi_j = \psi_j^T K \psi_i$) and M ($\psi_i^T M \psi_j = \psi_j^T M \psi_i$)

Mode shapes orthogonality

$$\psi_j^T M \psi_i = 0 \qquad i \neq j$$

Define $\mu_i = \psi_i^T M \psi_i \longrightarrow \psi_i^T M \psi_j = \delta_{ij} \mu_i$
 $K \psi_i = \omega_i^2 M \psi_i \longrightarrow \psi_i^T K \psi_j = \delta_{ij} \mu_i \omega_i^2$

Matrix notation

$$\Psi = \begin{bmatrix} \psi_1 & \psi_2 & \dots & \psi_n \end{bmatrix}$$

$$\Psi^T M \Psi = diag(\mu_i)$$

$$\Psi^T K \Psi = diag(\mu_i \omega_i^2)$$

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10

Example of a 2 DOFs system

$$\mathbf{x}_{\mathbf{x}_{1}} \quad \mathbf{x}_{\mathbf{x}_{2}} \quad \begin{bmatrix} m & 0 \\ 0 & m \end{bmatrix} \left\{ \begin{array}{c} \ddot{x}_{1} \\ \ddot{x}_{2} \end{array} \right\} + \begin{bmatrix} 2k & -k \\ -k & 2k \end{bmatrix} \left\{ \begin{array}{c} x_{1} \\ x_{2} \end{array} \right\} = \left\{ \begin{array}{c} 0 \\ f \end{array} \right\}$$

$$(K - \omega^{2}M) \psi = 0$$

$$det(K - \omega^{2}M) = det \begin{pmatrix} 2k - \omega^{2}m & -k \\ -k & 2k - \omega^{2}m \end{pmatrix} = 0$$

$$(2k - \omega^{2}m)(2k - \omega^{2}m) - k^{2} = m^{2}\omega^{4} - 4km\omega^{2} + 3k^{2} = 0$$
Second order equation in $\omega^{2} \quad \omega_{1}^{2} = k/m$

$$\omega_{2}^{2} = 3k/m$$

11

Example of a 2 DOFs system

$$(K - \omega^2 M) \psi = 0 \qquad (K - \omega^2 M) = \begin{bmatrix} 2k - \omega^2 m & -k \\ -k & 2k - \omega^2 m \end{bmatrix}$$
$$\psi = \begin{cases} A_1 \\ A_2 \end{cases}$$
 For $\omega^2 = k/m$

For $\omega_1^2 = k/m$

$$\left(2k - \frac{k}{m}m\right)A_1 - kA_2 = 0$$
$$kA_1 = kA_2 \Rightarrow A_1 = A_2$$

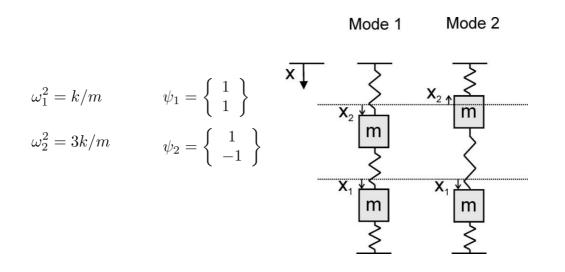
For $\omega_2^2 = 3k/m$

$$\left(2k - \frac{3k}{m}m\right)A_1 - kA_2 = 0$$
$$-kA_1 = kA_2 \Rightarrow A_1 = -A_2$$

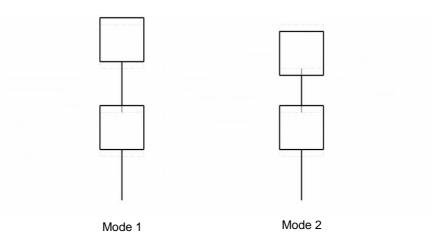
12

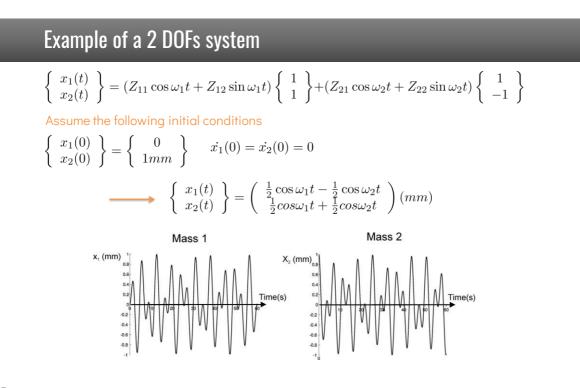
11

Example of a 2 DOFs system



Example of a 2 DOFs system





Resonance of MDOF systems



https://youtu.be/OaXSmPgl1os

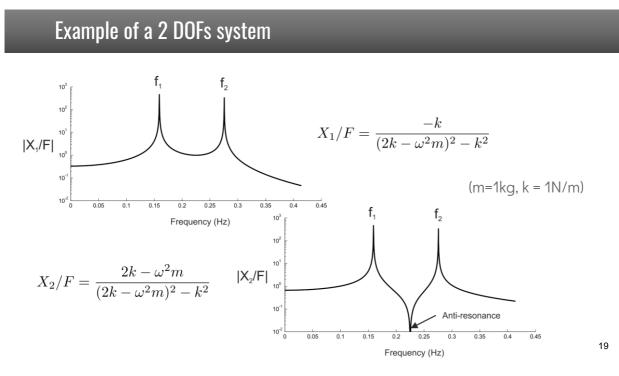
Harmonic excitation

$$\mathbf{x}_{\mathbf{x}_{1}} \underbrace{\mathbf{x}_{2}}_{\mathbf{x}_{2}} \mathbf{x}_{2} \underbrace{\mathbf{x}_{2}}_{\mathbf{x}_{2}} \mathbf{x}_{2} \underbrace{\mathbf{x}_{2}}_{\mathbf{x}_{2}} \left[\begin{array}{c} m & 0 \\ 0 & m \end{array} \right] \left\{ \begin{array}{c} \ddot{x}_{1} \\ \ddot{x}_{2} \end{array} \right\} + \left[\begin{array}{c} 2k & -k \\ -k & 2k \end{array} \right] \left\{ \begin{array}{c} x_{1} \\ x_{2} \end{array} \right\} = \left\{ \begin{array}{c} 0 \\ f \end{array} \right\}$$

$$\mathbf{x}_{1} \underbrace{\mathbf{x}_{2}}_{\mathbf{x}_{2}} \mathbf{x}_{2} \underbrace{\mathbf{x}_{1}}_{\mathbf{x}_{2}} \underbrace$$

Example of a 2 DOFs system

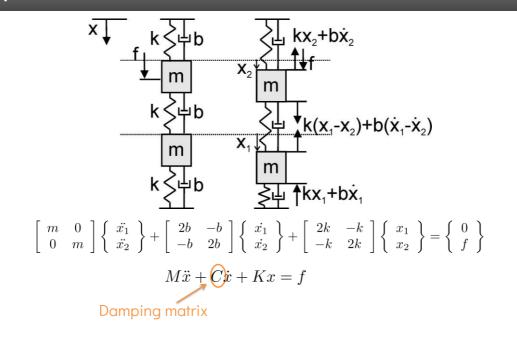
Response to harmonic excitation	$x_1(t) = X_1 e^{i\omega t}$ $x_2(t) = X_2 e^{i\omega t}$	$f(t) = F e^{i\omega t}$
$\left[\begin{array}{cc} 2k - \omega^2 m & -k \\ -k & 2k - \omega^2 m \end{array}\right] \left\{\begin{array}{c} X_1 \\ X_2 \end{array}\right\}$	$= \left\{ \begin{array}{c} 0\\ F \end{array} \right\}$	
$X_1/F = \frac{k}{(2k - \omega^2 m)^2 - k^2}$	$\omega_1^2 = k/m$ $\longrightarrow \omega_2^2 = 3k/m$	Resonance







Equations of motion



21

Free response

$$\left\{\begin{array}{c} x_1(t) \\ x_2(t) \end{array}\right\} = \left\{\begin{array}{c} A_1 \\ A_2 \end{array}\right\} e^{rt} = \psi e^{rt}$$

$$\left(K + rC + r^2M\right)\psi = 0$$

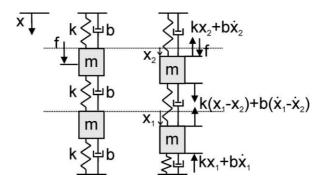
Non trivial solution if

 $det(K + rC + r^2M) = 0$

•Complex roots of the characteristic equation -> Oscillatory functions with exponential envelope

•Complex eigen vectors = complex modeshapes -> Not often used in practice in vibrations

Harmonic excitation



 $M\ddot{x} + C\dot{x} + Kx = f \qquad x(t) = Xe^{i\omega t} \qquad f(t) = Fe^{i\omega t}$

$$\left(K + i\omega C - \omega^2 M\right) X = F$$

23

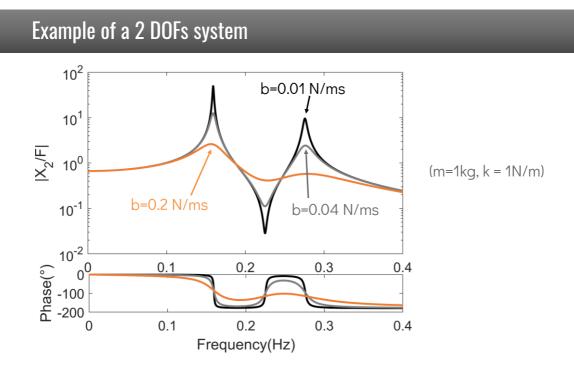
23

Example of a 2 DOFs system

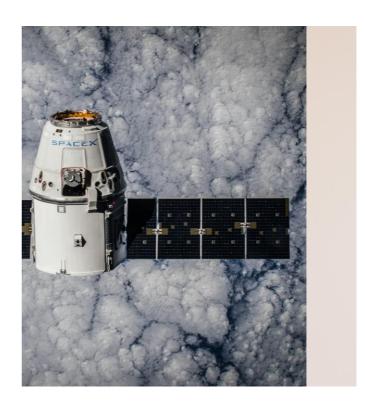
Response to harmonic excitation $x_1(t) = X_1 e^{i\omega t}$ $f(t) = F e^{i\omega t}$ $x_2(t) = X_2 e^{i\omega t}$

$$\begin{bmatrix} 2k+2i\omega b-\omega^2 m & -(k+i\omega b)\\ -(k+i\omega b) & 2k+2i\omega b-\omega^2 m \end{bmatrix} \begin{cases} X_1\\ X_2 \end{cases} = \begin{cases} 0\\ F \end{cases}$$

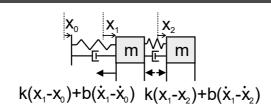
$$X_1/F = \frac{k + i\omega b}{(2k + 2i\omega b - \omega^2 m)^2 - (k + i\omega b)^2} \qquad \qquad \text{Damped resonances}$$
$$X_2/F = \frac{2k + 2i\omega b - \omega^2 m}{(2k + 2i\omega b - \omega^2 m)^2 - (k + i\omega b)^2} \qquad \qquad \text{No strict anti-resonance}$$







Base excitation of MDOF systems



Equations of motion:

 $m\ddot{x_1} = -k(x_1 - x_0) - b(\dot{x_1} - \dot{x_0}) - k(x_1 - x_2) - b(\dot{x_1} - \dot{x_2})$ $m\ddot{x_2} = k(x_1 - x_2) + b(\dot{x_1} - \dot{x_2})$

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27

Base excitation of MDOF systems

$$\left[\begin{array}{cc}m&0\\0&m\end{array}\right]\left\{\begin{array}{cc}\dot{x_{1r}}\\\dot{x_{2r}}\end{array}\right\}+\left[\begin{array}{cc}2b&-b\\-b&b\end{array}\right]\left\{\begin{array}{cc}\dot{x_{1r}}\\\dot{x_{2r}}\end{array}\right\}+\left[\begin{array}{cc}2k&-k\\-k&k\end{array}\right]\left\{\begin{array}{cc}x_{1r}\\x_{2r}\end{array}\right\}=\left\{\begin{array}{cc}-m\ddot{x_0}\\-m\ddot{x_0}\end{array}\right\}$$

Matrix notations:

$$M\ddot{x_r} + C\dot{x_r} + Kx_r = -M\ddot{x_b}$$

$$x_r = \left\{ \begin{array}{c} x_1 - x_0 \\ x_2 - x_0 \end{array} \right\} \qquad \qquad \ddot{x_b} = \left\{ \begin{array}{c} \ddot{x_0} \\ \ddot{x_0} \end{array} \right\} = \left[\begin{array}{c} 1 \\ 1 \end{array} \right] \ddot{x_0} = T \ddot{x_0}$$

All developments for force excitation apply